The role of time-dependent injection in blazar flares

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Figure: SEDs of PKS 2155-304 during 1 night in July 2006 (Aharonian et al., 2009, A&A 502, 749)

- Continuous injection of particles, changing other parameters
- Advantage: Solution of kinetic equation is well known, since it is in equilibrium
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What happens if the particle injection is time-dependent?

Kinetic equation for electrons:

$$\frac{\partial \textit{n}(\gamma, t)}{\partial t} - \frac{\partial}{\partial \gamma} \left[|\dot{\gamma}|_{\rm tot} \textit{n}(\gamma, t) \right] = \textit{Q}(\gamma, t)$$

Injection: one-time, relativistic

$$Q(\gamma, t) = q_0 \delta(\gamma - \gamma_0) \delta(t)$$

$$\begin{split} \dot{\gamma}|_{\rm tot} &= |\dot{\gamma}|_{\rm syn} + |\dot{\gamma}|_{\rm ec} + |\dot{\gamma}|_{\rm ssc} \\ &= D_0 (1 + l_{\rm ec}) \gamma^2 + A_0 \gamma^2 \int_0^\infty \, \mathrm{d}\gamma \, \gamma^2 n(\gamma, t) \end{split}$$

- Synchrotron cooling and external Compton cooling are linear and not time-dependent
- Synchrotron-self Compton cooling depends on the synchrotron field, which depends on the electron distribution ⇒ nonlinear (if injection time-dependent)
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Defined as the root of the ratio of the nonlinear to the linear cooling at t = 0

$$\alpha = \left[\frac{|\dot{\gamma}(t=0)|_{\rm ssc}}{|\dot{\gamma}|_{\rm syn} + |\dot{\gamma}|_{\rm ec}}\right]^{1/2} = \gamma_0 \left(\frac{q_0 A_0}{D_0(1+l_{ec})}\right)^{1/2}$$

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- The higher the electron density q₀ in the source, the higher the probability of (initial) nonlinear cooling

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Calculating the intensity:

$$I_i(
u,t) \propto \int \,\mathrm{d}\gamma \; n(\gamma,t) p_i(
u,\gamma)$$

• p_i: power of the respective emission process

Calculating the time-integrated SED

$$f_i(
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- Total SED: Unlimited integration $(0 \rightarrow \infty)$
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Total SED



- $\bullet\,$ Compton dominance is determined by α and $\mathit{I_{ec}}$
- $\bullet\,$ Total SED for $\alpha \ll 1$ resembles the SED of continuous injection
- In case of $\alpha \gg 1$ all components exhibit an additional break (position depends on α)
- Form of energy injection does not influence the main results (only the highest energies)

- \bullet Energy losses set in much earlier for $\alpha\gg 1$
- After t_c the SEDs become indistinguishable
- Time of observation is very important for modelling!

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- Zacharias & Schlickeiser 2012a, MNRAS 420, 84
- Zacharias & Schlickeiser 2010, A&A 524, A31
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- Zacharias & Schlickeiser 2013, ApJ 777, 109
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Thank you for your attention!

Backup

Light curves

Model:

- Complementary tool to obtain/check source parameters
- Using a simple model:

• Only retardation and geometry of the source taken into account:

$$L_{l}(\epsilon, t) = \int_{0}^{1} l_{i}(\epsilon, t - \lambda_{0}\lambda) \operatorname{H}\left[t - \lambda_{0}\lambda\right] \frac{dV(\lambda)}{V(\lambda)}$$

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- In case of $lpha \gg 1$ the variability time is reduced
- Light curves exhibit different temporal behaviour
- Shape of the light curves depends on the frequency



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Time-dependent injection in blazars



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