# The origin and evolution of cosmic magnetism

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## Abstract

Magnetism is one of the four fundamental forces. However, the origin of magnetic fields in stars, galaxies and clusters is an open problem in astrophysics and fundamental physics. When and how were the first fields generated? Are present-day magnetic fields a result of dynamo action, or do they represent persistent primordial magnetism? What role do magnetic fields play in turbulence, cosmic ray acceleration and galaxy formation? Here, we demonstrate how the Square Kilometre Array can deliver new data which will directly address these currently unanswered issues. Much of what we present is based on an *all-sky survey of rotation measures*, in which Faraday rotation towards $>10^7$ background sources will provide a dense grid for probing magnetism in the Milky Way, in nearby galaxies, and in distant galaxies, clusters and protogalaxies. Using these data, we can map out the evolution of magnetised structures from redshifts $z > 3$ to the present, can distinguish between different origins for seed magnetic fields in galaxies, and can develop a detailed model of the magnetic field geometry of the intergalactic medium and of the overall Universe. With the unprecedented capabilities of the SKA, the window to the Magnetic Universe can finally be opened.
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## 1. Introduction

Understanding the Universe is impossible without understanding magnetic fields. They fill intracluster and interstellar space, affect the evolution of galaxies and galaxy clusters, contribute significantly to the total pressure of interstellar gas, are essential for the onset of star formation, and control the density and distribution of cosmic rays in the interstellar medium (ISM). But in spite of their importance, the *evolution*, *structure* and *origin* of magnetic fields are all still open problems in fundamental physics and astrophysics. Specifically, we still do not know how magnetic fields are generated and maintained, how magnetic fields evolve as galaxies evolve, what the strength and structure of the magnetic field of the intergalactic medium (IGM) might be, or whether fields in galaxies and clusters are primordial or generated at later epochs.

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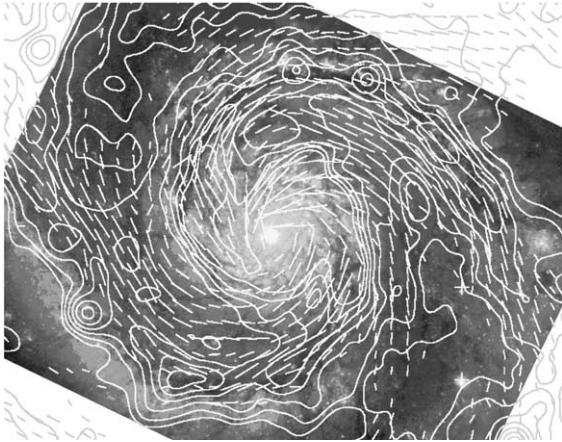*E-mail address:* bgaensler@cfa.harvard.edu (B.M. Gaensler).

Fig. 1. The magnetic field of the grand design spiral galaxy M 51. The image shows an optical HST image, overlaid with contours showing the radio total intensity emission at 5 GHz. The vectors show the orientation of the magnetic field, as determined from 5 GHz linear polarization measurements (Faraday rotation is small at this frequency) (Fletcher et al., 2004).
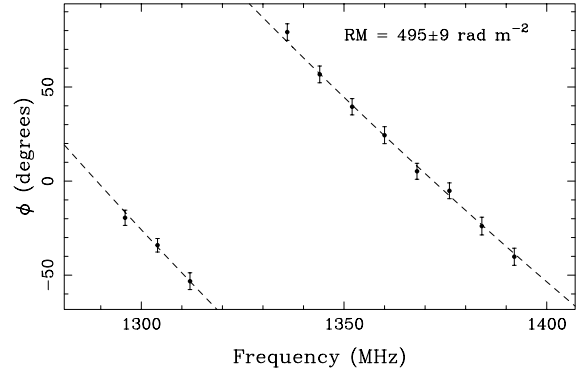


Fig. 2. Spectropolarimetry of pulsar B1154-62, demonstrating how the RM can be accurately determined using the position angle of polarization in multiple channels within the same observing band (Gaensler et al., 1998).

Ultimately, we would like to establish whether there is a connection between magnetic field formation and structure formation in the early Universe, and to obtain constraints on when and how the first magnetic fields in the Universe were generated.

Most of what we know about astrophysical magnetic fields comes through the detection of radio waves. *Synchrotron emission* measures the total field strength, while its *polarization* yields the orientation of the regular field in the sky plane and also gives the field's degree of ordering (see Fig. 1). Incorporating *Faraday rotation* yields a full three-dimensional view by providing information on the field component along the line of sight (see Fig. 2), while the *Zeeman effect* provides an independent measure of field strength in cold gas clouds. However, measuring astrophysical magnetic fields is a difficult topic still in its infancy, restricted to nearby or bright objects. In this chapter we explain how the Square Kilometre Array (SKA) can revolutionise the study of cosmic magnetism.

## 2. Polarimetry with the SKA

Much of what the SKA can contribute to our understanding of magnetic fields will come from

its *polarimetric capabilities*. As discussed in more detail elsewhere (Beck and Gaensler, 2004; Feretti, 2004), the main SKA specifications which enable this capability will be *high polarization purity* and *spectropolarimetric capability*. The former will allow detection of the relatively low linearly polarized fractions ($\lesssim 1\%$) from most astrophysical sources seen down to sub-mJy levels, while the latter will enable accurate measurements of Faraday rotation measures (RMs), intrinsic polarization position angles and Zeeman splitting. Other minimum requirements include a frequency coverage of 0.5–10 GHz, a field of view at 1.4 GHz of 1 deg$^2$ which can be fully imaged at 1″ resolution, a significant ($\sim 50\%$) concentration of the collecting area into a central core of diameter $\sim 5$ km, and a longest baseline of $\sim 300$ km.

It is important to appreciate how powerful the combination of high sensitivity, good polarization purity and full spectropolarimetry of the SKA will be. With current instruments, the only way to simultaneously determine accurate values for both the RM and intrinsic polarization position angle is to make many observations across a broad frequency range. This is not only time-consuming, but the analysis must proceed cautiously, since various depolarizing effects have a strong frequency dependence (e.g., Sokoloff et al., 1998).

The high sensitivity and broad bandwidth of the SKA eliminate these difficulties: a single spectro-

polarimetric observation at a single IF can simultaneously provide good estimates of both RM and position angle, the limiting factor often only being the accuracy of ionospheric corrections to the observed Faraday rotation. For example, at an observing frequency of 1.4 GHz with a fractional bandwidth of 25%, a 1-min SKA observation of a source with a linearly polarized surface brightness of $\sim 8$ μJy beam$^{-1}$ will yield a RM determined to an accuracy $\Delta RM \approx \pm 5$ rad m$^{-2}$, and its intrinsic position angle measured to within $\Delta \Theta \approx \pm 10°$. Measurements of this precision will be routinely available in virtually any SKA observation of a polarized source.

## 3. All-sky rotation measures

Currently $\sim 1200$ extragalactic sources and $\sim 300$ pulsars have measured RMs. These data have proved useful probes of magnetic fields in the Milky Way, in nearby galaxies, in clusters, and in distant Lyα absorbers. However, the sampling of such measurements over the sky is very sparse, and most measurements are at high Galactic latitudes.

A key platform on which to base the SKAs studies of cosmic magnetism will be to carry out an *all-sky RM Survey*, in which spectropolarimetric continuum imaging of 10000 deg$^2$ of the sky can yield RMs for approximately $2 \times 10^4$ pulsars and $2 \times 10^7$ compact polarized extragalactic sources in about a year of observing time (see Beck and Gaensler, 2004 for a detailed description). This data set will provide a grid of RMs at a mean spacing of $\sim 30'$ between pulsars and just $\sim 90''$ between extragalactic sources (see Fig. 3).

In the following sections, we describe some of the SKA experiments which will provide new insight into the origin, structure and evolution of cosmic magnetic fields. Many of these observations stem directly from the RM grid just described, or from deeper observations of specific regions which will provide an even more closely spaced set of background RMs. This discussion begins with the Milky Way and its ISM, and then considers other sources at increasingly larger distances.
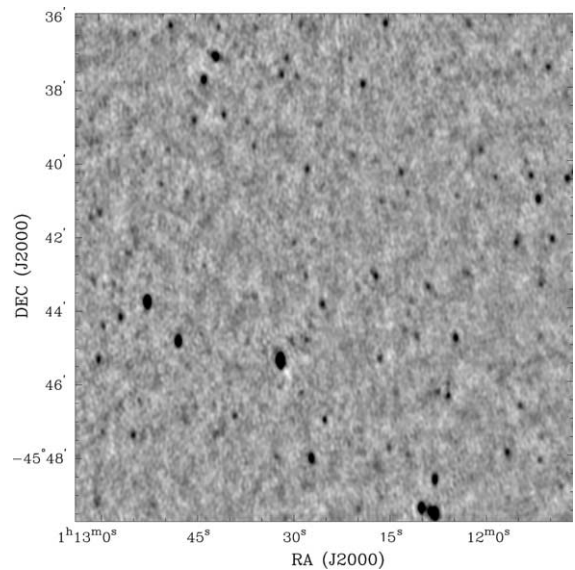


Fig. 3. 1.4 GHz ATCA image of a small section of the Phoenix Deep Field (Hopkins et al., 2003). The density of extragalactic sources in this total intensity image is similar to what will be achievable in *polarized intensity* in a 1-h 1.4 GHz observation with the SKA.

## 4. The local universe

### 4.1. The milky way and nearby galaxies

Key information on the processes which amplify and sustain large-scale magnetic field structure in galaxies comes from characterising and comparing the three-dimensional structure of the magnetic field in nearby galaxies. Amongst the questions which we would like to resolve are the number and location of magnetic reversals, the relation between magnetic and optical spiral arms, the vertical magnetic structure in the disk, the strength and structure of magnetic fields in galactic halos, and the correspondence between magnetic field structure and other parameters such as rotation speed, Hubble type and star formation rate.

In our Milky Way, the large sample of pulsar RMs obtained with the SKA, combined with distance estimates to these sources from parallax or from their dispersion measures, can be inverted to yield a complete delineation of the magnetic field in the spiral arms and disk on scales $\gtrsim 100$ pc (e.g., Stepanov et al., 2002). Small-scale

structure and turbulence can be probed using *Faraday tomography*, in which foreground ionised gas produces complicated frequency dependent Faraday features when viewed against diffuse Galactic polarized radio emission (e.g., Gaensler et al., 2001; Haverkorn et al., 2003; Uyanıker, 2003). Magnetic field geometries in the Galactic halo and outer parts of the disk can be studied using the extragalactic RM grid (see detailed discussion by Beck and Gaensler, 2004).

In external galaxies, magnetic fields can be directly traced by diffuse synchrotron emission and its polarization. For the nearest galaxies, the SKA will provide sufficient resolution and sensitivity to carry out Faraday tomography on this emission, allowing us to identify individual features, structures and turbulent processes in the magneto-ionised ISM of these galaxies. These measurements can be compared with Hɪ, Hα and continuum observations at a variety of wavelengths to establish, in a much more direct way than is possible for the confused lines of sight of our own Galaxy, how magnetic features interact with other features in the ISM.

Synchrotron emission is seen only in magnetised regions in which either cosmic ray electrons have been accelerated, or into which such electrons have diffused. In other parts of a galaxy, particularly in the outer disk and halo, we must rely on background RMs to map magnetic fields. However, the current density of such sources is such that only a handful of nearby galaxies can be studied in this way (e.g., Han et al., 1998). With the sensitivity of the SKA, deep observations of nearby galaxies can provide $>10^5$ background RMs, and thus allow fantastically detailed maps of the magnetic structure. Even at a distance of 10 Mpc, most galaxies should be traversed by $\sim$50 sightlines to polarized background sources. The sample of galaxies which could be studied in this way is very large.

As discussed in detail by Beck and Gaensler (2004), the overall structure of galactic magnetic fields provides a potential discriminant between dynamo and primordial mechanisms for the origin and sustainment of these fields (Beck et al., 1996). For example, the mean-field α–Ω dynamo model makes specific predictions for the symmetries seen in both azimuthal and vertical field structure. Observations of background RMs and of diffuse polarized synchrotron emission can allow direct measurements of the azimuthal modes, but these measurements are limited at present to $\sim$20 galaxies (Beck, 2000). With the SKA, this sample can be increased by up to three orders of magnitude. These data can allow us to distinguish between different conditions (e.g., strong density waves, high star formation rate, or interactions) for excitation of various dynamo modes. Meanwhile, the presence and prevalence of reversals in the disk field structure, plus the structure of field in the halo, will together let us distinguish between dynamo and primordial models for field origin in different systems.

## 4.2. Galaxy clusters

In clusters of galaxies, magnetic fields play a critical role in regulating heat conduction (Chandran and Cowley, 1998; Narayan and Medvedev, 2001), and may also both govern and trace cluster formation and evolution. Estimates of the overall magnetic field strength come from inverse compton detections in X-rays, from detection of diffuse synchrotron emission, from cold fronts and from simulations, but our only direct measurements of field strengths and geometries come from RMs of background sources (see Feretti, 2004). Currently just $\sim$1–5 such RM measurements can be made per cluster (e.g., Govoni et al., 2001); only by considering an ensemble of RMs averaged over many systems can a crude picture of cluster magnetic field structures be established (Clarke et al., 2001).

With the SKA, the RM grid can provide $\sim$1000 background RMs behind a typical nearby cluster; a comparable number of RMs can be obtained for a more distant cluster through a deep targeted observation. These data will allow us to derive a detailed map of the field in *each* cluster. With such information, careful comparisons of the field properties in various types of cluster (e.g., those containing cooling flows, those showing recent merger activity, etc.) at various distances can be easily made. Furthermore, detailed comparisons between RM distributions and X-ray images of clusters will become possible, allowing us to relate

the efficiency of thermal conduction to the magnetic properties of different regions, and to directly study the interplay between magnetic fields and hot gas.

Deep observations targeting individual clusters will also allow the detection of low surface brightness synchrotron emission, allowing us to explore the role of low-level magnetic fields, and providing vital clues to the evolution of cluster magnetic fields.

## 5. Galaxies at intermediate redshifts

Many galaxies at intermediate redshifts $(0.1 \lesssim z \lesssim 2)$, are representative of the local population but at earlier epochs. Measurements of the magnetic field in such systems thus provides direct information on how magnetised structures evolve and amplify as galaxies mature. However, the linearly polarized emission from galaxies at these distances will often be too faint to detect directly; Faraday rotation thus holds the key to studying magnetism in these distant sources.

### 5.1. Galaxies in the foreground of bright extended radio sources

At intermediate redshifts, we expect galaxies to be <1′ in extent, too small to be adequately probed by the RM grid of compact background sources discussed in Section 3. However, there are many distant *extended* polarized sources (e.g., many of the quasars and radio galaxies in the 3C catalogue), which provide ideal background illumination for probing Faraday rotation in galaxies which happen to lie along the same line of sight. These experiments can deliver maps of magnetic field structures in galaxies more than 100 times more distant than discussed in Section 4.1, and thus provide information on the evolution of magnetic fields as a function of cosmic epoch. Currently this technique has been applicable in just a few fortuitous cases, such as for NGC 1310 seen projected against a lobe of Fornax A (Fig. 4; Fomalont et al., 1989; Schulman and Fomalont, 1992), and for an absorption line line system at $z = 0.395$ seen against PKS 1229–021 (Kronberg
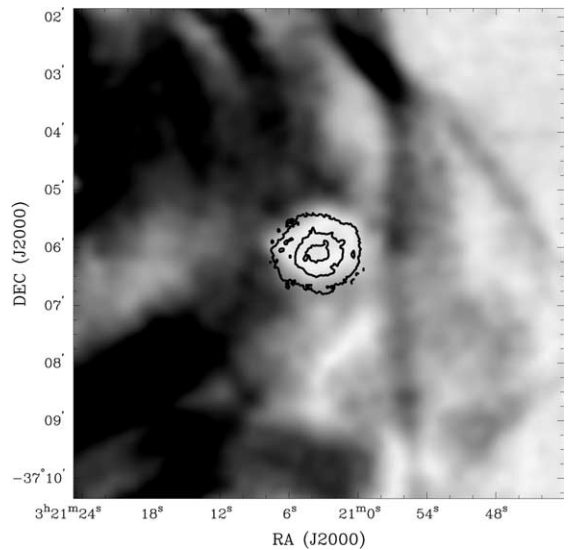
Fig. 4. A 1.4 GHz VLA image of linearly polarized intensity from a small region in the western lobe of Fornax A, superimposed with optical contours from the Digitized Sky Survey. The bright optical source corresponds to the spiral galaxy NGC 1310, which produces strong depolarization through its Faraday rotation towards the background polarized emission from Fornax A (Fomalont et al., 1989).

et al., 1992). With the greater sensitivity of the SKA, many such systems lying in front of bright polarized radio sources can be targeted. These objects are less evolved counterparts to the nearby galaxies discussed in Section 4.1; comparison of the two populations will demonstrate how present day magnetic fields emerge. For example, the $z = 0.395$ system discussed by Kronberg et al. (1992) appears to be a spiral galaxy with a bisymmetric magnetic field geometry and a magnetic field strength of $\sim$1–4 µG, surprisingly similar to results from nearby galaxies. Clearly a significant expansion of the sample of such systems can provide strong constraints on the processes which generate magnetic fields in galaxies.

### 5.2. RM statistics of Lyα absorbers

While the approach in Section 5.1 allows us to study a variety of individual systems in detail, an even larger sample, suitable for statistical studies, can be accumulated by considering unresolved

sources in the foreground of our RM grid. Specifically, we expect that a large number of the sources for which we measure RMs will be quasars showing foreground Lyα absorption; these absorption systems likely represent the progenitors of present-day galaxies. If a large enough sample of RMs for quasars at known redshift can be accumulated, a trend of RM vs *z* can potentially be identified. The form of this trend can then be used to distinguish between RMs resulting from magnetic fields in the quasars themselves and those produced by fields in foreground absorbing clouds (e.g., Welter et al., 1984); detection of the latter effect would then directly trace the evolution of magnetic field in galaxies and their progenitors. This experiment has been attempted several times with existing data sets, but only a relatively small number of absorbing systems show a clear excess in RM. These data correspondingly provide only marginal (if any) evidence for any evolution of RM with redshift (Welter et al., 1984; Wolfe et al., 1992; Perry et al., 1993; Oren and Wolfe, 1995). There are several limitations in these studies, all of which can be circumvented using the SKA.

First, all quasar RM measurements are contaminated by the Galactic contribution to the RM. The Galactic signal can be removed by considering the RMs of neighboring background sources, but currently the sampling of such nearby sources is extremely sparse, so that large uncertainties are introduced through this foreground subtraction. The SKA's RM grid will improve this situation by five orders of magnitude, allowing a complete template for the Galactic foreground to be derived and then removed (e.g., Frick et al., 2001).

Second, because *detected* RMs are lower than *intrinsic* RMs by a factor $(1 + z)^2$, where *z* is the redshift of the source in which Faraday rotation is produced, an absorber at a redshift of ∼3 must have an RM signal ∼80 times the measurement uncertainty in order to be considered significant. For many previous efforts, the uncertainties introduced by removing the Galactic foreground are as large as $\Delta RM \sim 20-30$ rad m$^{-2}$, making it difficult to detect magnetic fields in distant objects unless the intrinsic RM is extremely high. With the SKA accurate RM measurements and precise fore-

ground subtraction will allow us to reliably probe magnetic field evolution to much higher redshifts than has been previously possible.

Finally, an overriding limitation of previous studies has been the limited number of sources per redshift bin. By correlating SKA RMs with new surveys such as SDSS, a vast increase in the sample size is expected, so that different evolutionary models can be clearly distinguished.

## 6. Galaxies and protogalaxies at $z \gtrsim 2$

At yet higher redshifts, we can take advantage of the sensitivity of the deepest SKA fields, in which we expect to detect the synchrotron emission from the youngest galaxies and proto-galaxies (see contribution by Carilli, this volume). The tight radio-infrared correlation, which holds also for distant infrared-bright galaxies (Carilli et al., 2001), tells us that magnetic fields with strengths similar or even larger than in nearby galaxies existed in some young objects, but the origin of these fields is unknown. The total intensity of synchrotron emission can yield approximate estimates for the magnetic field strength in these galaxies. between magnetic fields and cosmic rays. In deep observations, the SKA can detect and resolve spiral galaxies like M 51 out to redshifts $z \approx 2$, and to even larger distances if the star formation rate is higher than in M 51 (see contribution by van der Hulst, this volume). Since a dynamo needs a few rotations or about $10^9$ year to build up a coherent field (Beck et al., 1994), detection of synchrotron emission in young galaxies at high *z* put constraints on the seed field which potentially challenge dynamo models (Perry et al., 1993).

Even very distant galaxies unresolvable by the SKA may reveal polarized emission if the inclination is high or if the field structure is asymmetric, e.g., due to tidal interaction or ram pressure as has been observed in several galaxies (e.g., in NGC 2276; Hummel and Beck, 1995). Interactions were much more frequent in the early Universe so excess polarization might be expected in some objects, especially if magnetic fields were already present before compression.

## 7. The IGM and cosmic field geometry

Fundamental to all the issues discussed above is the search for magnetic fields in the IGM. All of "empty" space may be magnetised, either by outflows from galaxies, by relic lobes of radio galaxies, or as part of the "cosmic web" structure. Such a field has not yet been detected, but its role as the likely seed field for galaxies and clusters, plus the prospect that the IGM field might trace and regulate structure formation in the early Universe, places considerable importance on its discovery. Once the magnetic field of the IGM is detected, a measurement of the characteristic size scales of its fluctuations can allow us to differentiate between the wide variety of mechanisms proposed for magnetic field generation in the IGM (see Kronberg, 1994; Grasso and Rubinstein, 2001; Carilli and Taylor, 2002; Widrow, 2002; Giovannini, 2003 for extensive reviews). While RMs of distant sources are potentially a useful probe of the IGM, to date there has been no detection of magnetic fields in the IGM; current upper limits on the strength of any such field are model dependent, but suggest $|B_{\text{IGM}}| \lesssim 10^{-8}$ to $10^{-9}$ G (Kronberg, 1994; Blasi et al., 1999).

Using the SKA, this all-pervading cosmic magnetic field can finally be identified through the RM grid proposed in Section 3. Just as the correlation function of galaxies yields the power spectrum of matter, the analogous correlation function of this RM distribution can provide the magnetic power spectrum of the IGM as a function of cosmic epoch and over a wide range of spatial scales. Such measurements will allow us to develop a detailed model of the magnetic field geometry of the IGM and of the overall Universe.

As mentioned earlier, the observed RM is reduced by a factor $(1 + z)^2$ over its intrinsic value. However, a simplistic assumption is that the comoving magnetic flux and electron density are both constant, so that electron density evolves as $(1 + z)^3$, while the magnetic field evolves as $(1 + z)^2$ (e.g., Widrow, 2002). In this case, we expect the observed RM to be proportional to $(1 + z)^3$. While this is a naive calculation, it demonstrates the general principle that RM signatures from distant objects can be quite large, even if the
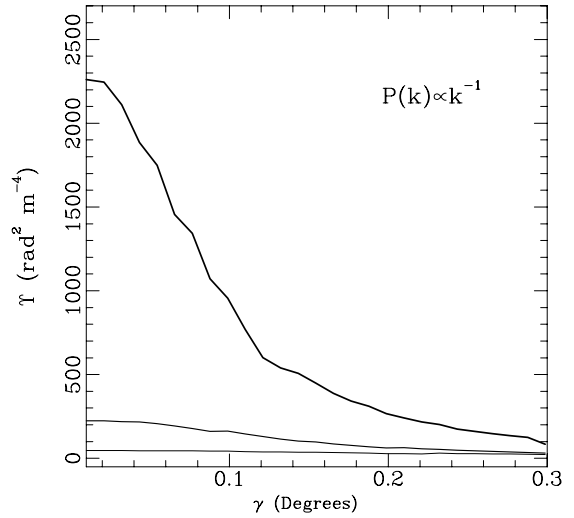


Fig. 5. RM correlation between pairs of background sources as a function of separation angle, assuming a power spectrum for magnetic fluctuations of the form $P_{\text{B}}(k) \propto k$ (Kolatt, 1998). From bottom to top, the three lines show the correlation function for sources at redshifts of $z = 0.5$, 1 and 2. Reproduced by permission of the AAS.

foreground IGM has a relatively small magnetic field.

More robust treatments are presented by Kolatt (1998) and by Blasi et al. (1999), who demonstrate in detail how the statistics of RM measurements as a function of redshift can be used to remove the foreground Galactic contribution, and then to extract the strength and power spectrum of magnetic fields in the IGM. These studies require a large sample of RMs from sources at known redshift, which should be obtainable by combining the SKA RM grid with the wide-field spectroscopic data bases provided by SDSS and KAOS, plus with the all-sky multiband photometry of LSST and SkyMapper. For example, Kolatt (1998) shows that an IGM field with strength $10^{-10}$ G can be detected with 5-$\sigma$ significance by considering extragalactic RM measurements over just a few degrees of sky, provided that the density of RM measurements is $\sim$800 sources deg$^{-2}$; the resulting correlation function is shown in Fig. 5. Figure 2 of Beck and Gaensler (2004) demonstrates that this is well beyond the capabilities of current telescopes, but is easily achievable with
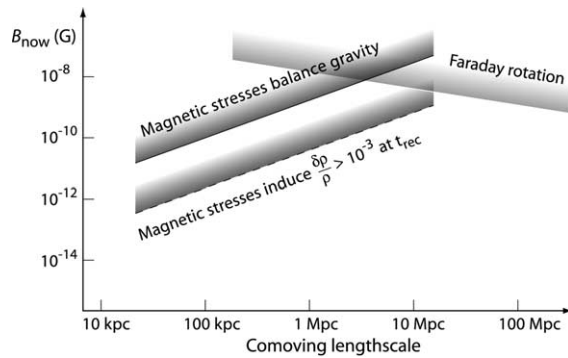
Fig. 6. Observational and theoretical constraints on the magnetic field on various length scales at the time of recombination, $t_{rec}$. A field is cosmogonically important if it can generate density perturbations of amplitude $\sim 10^{-3}$ at $t_{rec}$, since these would have developed into gravitationally bound systems by the present time (Rees, 2000).

the SKA. We note that this experiment requires sub-arcsec spatial resolution at radio wavelengths, so as to ensure accurate optical identification of polarized radio sources.

## 8. Primordial fields?

As emphasised elsewhere in this volume, a large part of SKA science will be unanticipated and serendipitous discoveries, which do not easily extrapolate from previous work. Thus, although speculative, we suggest here some isolated experiments which might provide some constraints on the strength and origin of magnetic fields at very early epochs.

A magnetic field already present at the recombination era might have affected the cosmogonic process (Subramanian and Barrow, 1998). The constraints are summarised in Fig. 6, in terms of the field's characteristic length scale.

A variety of cosmological processes can produce relatively high primordial magnetic fields, of strength $10^{-10}$–$10^{-9}$ G at $z \sim 5$ (see Grasso and Rubinstein, 2001). Active galactic nuclei and violent star-formation activity in young galaxies may be able to generate seed fields of similar strength (e.g., Chakrabarti et al., 1994; Kronberg et al., 1999; Völk and Atoyan, 2000;
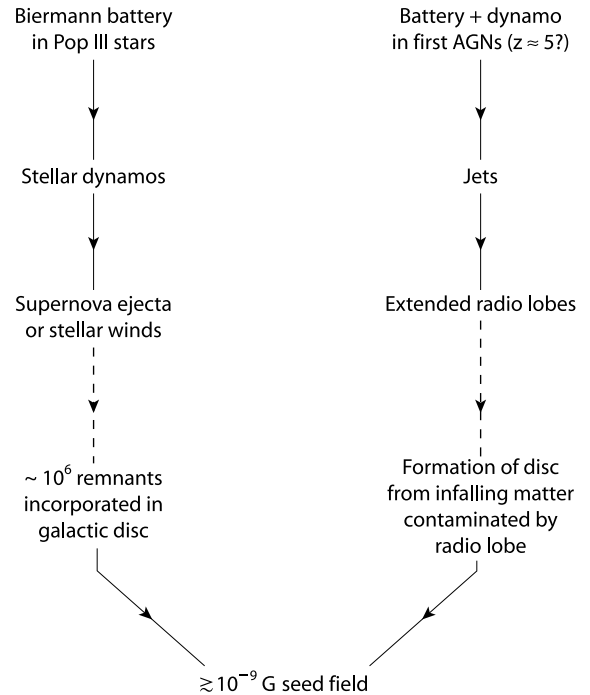


Fig. 7. Suggested mechanisms through which Population III stars and active galactic nuclei can both produce seed magnetic fields of strength $\gtrsim 10^{-9}$ G (Rees, 2004).

Furlanetto and Loeb, 2001), as illustrated schematically in Fig. 7. The dynamo mechanism may then have amplified the seed field to the microgauss strength observed in galaxies today. RM experiments with the SKA toward very high redshift sources might provide evidence for such fields, taking advantage of the RM $\propto (1 + z)^3$ effect discussed in Section 7. For example, we already know of $\gamma$-ray bursts at redshifts as high as $z = 4.5$ (Andersen et al., 2000), and of radio galaxies at $z = 5.2$ (van Breugel et al., 1999). With the advent of projects such as *SWIFT* and LOFAR, it is virtually certain that the sample of such high redshift sources will soon greatly expand. With the high sensitivity of the SKA, linear polarization and Faraday rotation should be detectable from some of these sources. Using a foreground template at lower redshift provided by the RM grid, the RM contribution at high redshifts can then be isolated, yielding a possible detection of magnetic fields at early epochs.

## 9. Conclusion

We have outlined the exciting new insights which the SKA can provide into the origin, evolution and structure of cosmic magnetic fields. The sheer weight of RM statistics which the SKA can accumulate, combined with deep polarimetric observations of individual sources, will allow us to characterise the geometry and evolution of magnetic fields in galaxies, in clusters and in the IGM from high redshifts through to the present. We may also be able to provide the first constraints on when and how the first magnetic fields in the Universe were generated. Apart from these experiments which we can conceive today, we also expect that the SKA will certainly discover new magnetic phenomena beyond what we can currently predict or even imagine.

Clearly this set of experiments meets all the criteria for SKA Key science: it addresses several unanswered questions in fundamental astrophysics, is science which exploits the unique polarimetric capabilities of radio astronomy and of the SKA, focuses on a topic ("magnetism") which is easily accessible to the general public, and is an important part of the themes of galaxy formation and cosmic evolution, currently considered high priority by funding agencies.

The main SKA specifications which enable us to reach our goals are high polarization purity, spectropolarimetric capability, and a frequency coverage of at least 0.5–10 GHz. The all-sky RM Survey is best done at a frequency of 1.4 GHz, perhaps simultaneously with other wide-field continuum and H I surveys. Faraday tomography requires frequencies $\sim$0.5–1 GHz for low-density, and $\sim$1–5 GHz for high-density regions of the local ISM. Mapping of the polarized synchrotron emission in the Milky Way, external galaxies and clusters is best done at frequencies $\gtrsim$ 5–10 GHz where Faraday depolarization is minimal. Other minimum requirements are a field of view at 1.4 GHz of 1 deg$^2$ which can be fully imaged at 1$''$ resolution, a significant ($\sim$50%) concentration of the collecting area into a central core of diameter $\sim$5 km, and a largest baseline of $\sim$300 km.

To summarise, incredibly rich data sets for the study of magnetic fields await the advent of the SKA. With the unique sensitivity, resolution and polarimetric capabilities of this next-generation radio telescope, we can finally address fundamental questions relating to the physics of cosmic magnetic fields, and their role in forming structure on all scales.

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